

Brane-Higgs fields

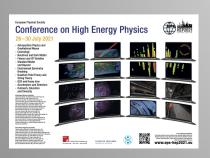


Grégory Moreau Pôle de Physique Théorique

R.Leng, GM, F.Nortier, PRD 103 (2021) 075010

A.Angelescu, R.Leng, GM, F.Nortier, PRD 101 (2020) 075048





Outline

- A The scenarii
- **B** The methodologies
 - **C** Beyond Higgs regularisation
 - **D** Wave function jumps
 - **E UV origin of chirality**
- F Phenomenological impacts

A - The scenarii

Framework:

Higgs boson at a point along warped extra dimension(s)

hep-ph/9905221

[3-brane] where gravity scale is reduced down to TeV!

Randall, Sundrum

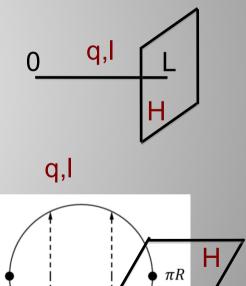
=> no more gauge hierarchy problem (with SM scale)

I) Interval model

- Fermions in the bulk (for FCNC, flavours,...),
- Toy model with flat compact space.

II) S¹/Z₂ Orbifold model

$$egin{aligned} \mathcal{L}\left[\Phi(x^{\mu},-y)
ight] &= \mathcal{L}\left[\Phi(x^{\mu},y)
ight] \ \Phi(x^{\mu},-y) &= \mathcal{T}\Phi(x^{\mu},y) \end{aligned} egin{aligned} \mathcal{L}_{kin} &= rac{i}{2} \, \left(ar{Q}\Gamma^{M} \overleftrightarrow{\partial_{M}} Q + ar{D}\Gamma^{M} \overleftrightarrow{\partial_{M}} D
ight) \end{aligned}$$



B – The methodologies

5D approach:



All symmetries Field content Geometrical set-up



Action minimisation \longleftrightarrow Current condition

Natural BC & Essential BC

+ (with BC selection)

Equations of motion



Mass spectrum,

Profiles

$$\begin{split} \delta_{\bar{F}} S_{bulk} &= \int d^4 x \, \left(\int_{-\pi R^+}^{0^-} + \int_0^{\pi R} \right) dy \, \left\{ \delta \bar{F} \frac{\partial \mathcal{L}_{kin}}{\partial \bar{F}} + \delta \left(\partial_M \bar{F} \right) \frac{\partial \mathcal{L}_{kin}}{\partial \partial_M \bar{F}} \right\} \\ &= \int d^4 x \, \left(\int_{-\pi R^+}^{0^-} + \int_0^{\pi R} \right) dy \, \left\{ \delta \bar{F} \frac{\partial \mathcal{L}_{kin}}{\partial \bar{F}} + \partial_M \left[\delta \bar{F} \frac{\partial \mathcal{L}_{kin}}{\partial \partial_M \bar{F}} \right] - \delta \bar{F} \, \partial_M \frac{\partial \mathcal{L}_{kin}}{\partial \partial_M \bar{F}} \right\} \end{split}$$

$$Q_L\left(x^{\mu},y
ight) = rac{1}{\sqrt{2\pi R}} \sum_{n=0}^{+\infty} q_L^n(y) \, \psi_L^n\left(x^{\mu}
ight)$$

Barcelo, Mitra, GM arXiv:1408.1852 [hep-ph]

4D approach for KK tower masses with Yukawa couplings:

- 1) use free 5D method results (KK masses & profiles),
- 2) bi-diagonalise effective 4D field mass matrix (mixings).

C – Beyond Higgs regularisation

I) No Brane-Higgs regularisation (width to 0)

- Two regularisation processes non physically equivalent.
- No theoretical motivation, no guarantee to remain in same model.
- Mathematical inconsistencies like mixing functions and distributions.

=> irrelevant debate on ggF calculation non-commutativity

Carena; Neubert; Toharia; Goertz... initial paper: arXiv:1303.5702 [hep-ph]

II) EBC or BBT

- Essential Boundary Conditions (**EBC**, not Natural BC) are **necessary**: vanishing probability fermion currents [in both dual models].

- Their rôle can be played by new Bilinear Brane Terms (BBT).

$$S_{B} = \int d^{4}x \left(\sigma_{0}^{Q} \ \bar{Q}Q \Big|_{0} + \sigma_{\pi R}^{Q} \ \bar{Q}Q \Big|_{\pi R} + \sigma_{0}^{D} \ \bar{D}D \Big|_{0} + \sigma_{\pi R}^{D} \ \bar{D}D \Big|_{\pi R} \right)$$

...like in GR context : AdS/CFT duality, Gibbons-Hawking and scalar terms

III) Result overview

From mathematically rigorous analyses...

Free case

1)
$$(--): f_L^n(y) = B_L^n \sin(m_n y), (++): f_R^n(y) = B_L^n \cos(m_n y); \sin(m_n \pi R) = 0,$$

2)
$$(++)$$
: $f_L^n(y) = B_R^n \cos(m_n y)$, $(--)$: $f_R^n(y) = -B_R^n \sin(m_n y)$; $\sin(m_n \pi R) = 0$,

3)
$$(-+): f_L^n(y) = B_L^n \sin(m_n y), (+-): f_R^n(y) = B_L^n \cos(m_n y); \cos(m_n \pi R) = 0,$$

4)
$$(+-)$$
: $f_L^n(y) = B_R^n \cos(m_n y)$, $(-+)$: $f_R^n(y) = -B_R^n \sin(m_n y)$; $\cos(m_n \pi R) = 0$.

Yukawa coupling

$$\begin{cases} (+\times) : q_L^n(y) = A_q^n \cos(M_n y), & (-\times) : q_R^n(y) = -A_q^n \sin(M_n y), \\ (-\times) : d_L^n(y) = A_d^n \sin(M_n y), & (+\times) : d_R^n(y) = A_d^n \cos(M_n y), \end{cases}$$

$$\tan(M_n \pi R) = \left| \frac{X}{2} \right|, \ A_q^n = e^{i(\alpha_0^n + \alpha_Y)}, \ A_d^n = e^{i\alpha_0^n},$$
$$\tan(M_n \pi R) = -\left| \frac{X}{2} \right|, \ A_q^n = e^{i(\alpha_0^n + \alpha_Y \pm \pi)}, \ A_d^n = e^{i\alpha_0^n},$$

	No boundary characteristic	Vanishing current condition [EBC]	Bilinear brane terms [NBC]
4D Approach	(Impossible)	BC (±)	BC (±)
5D Approach	(Impossible)	(Impossible)	BC (×)

D – Wave function jumps

I) Interval models

No fermion profile discontinuities.

II) Orbifold models

Fermion profile jumps arise!

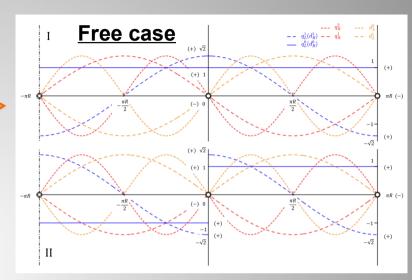
- Mathematically consistent,

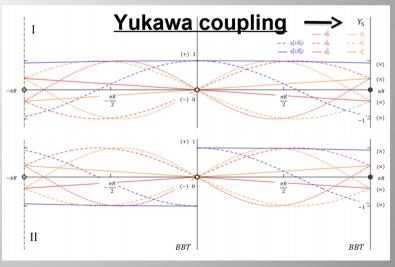
$$oxed{S_{bulk} = \int d^4x \left(\int_{-\pi R^+}^{0^-} dy \; \mathcal{L}_{kin} + \int_{0}^{\pi R} dy \; \mathcal{L}_{kin}
ight)} \; \; f(0) = f(0^+)}$$

- disappear for some free solution parities
 but unavoidable with brane-Yukawa couplings,
- physical impact neither on KK mass spetrum nor on 4D effective Yukawa couplings,

$$S_Y = \int d^4x \, {\cal L}_Y(x^\mu,\pi R) \, egin{aligned} {\cal L}_Y = -Y_5 \,\, H Q_L^\dagger D_R \end{aligned}$$

- models even with physical jumps probably exist...





 $-Y_5' HQ_B^{\dagger}D_L$

E – UV origin of chirality

The choice of EBC type (or equivalently of BBT) – via inclusive parity SYM. – generates the chiral nature of the low-energy model and the SM field chiralities.

F - Phenomenological impacts

No 'wrong-chirality' Yukawa coupling dependence (4D/5D method):

arXiv:0906.1990 [hep-ph]

Azatov et al.

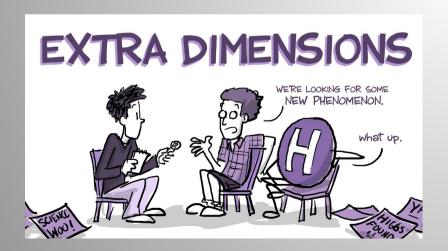
=> No significant FC quark Yukawa interactions from misalign. => no strong KK mass **constraints** from $\bar{K} - K$, $\bar{B} - B$

=> And no significant FC quark/lepton Yukawa couplings => no **detectable** exotic decays $t \to ch$ or $h \to \mu \tau$ (at LHC, LC) ?

Conclusions

Rigorous treatments of brane-Higgs scenarii:

- No brane-Higgs regularisation
- EBC or BBT: outside or inside the action
- Profile discontinuities via improper integrals
- Path towards UV origin of chirality





Only soften potential New Physics FC effects in the Higgs sector at LHC (LC)...